

transition metal analogues, Nb^v and Ta^v, do not. Iyer & Smith (1967) have noted similarly that there is no structural resemblance between the rare-earth double oxides of Pa^v and the analogous compounds of Nb^v and Ta^v.

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The Crystal Structure of 8,16-Oxido-*cis*-[2.2]metacyclophane

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(Received 27 November 1967)

The crystal structure of 8,16-oxido-*cis*-[2.2]metacyclophane, C₁₆H₁₄O, is orthorhombic, space group probably *Cmc*2₁, with $a = 8.78$, $b = 16.04$, $c = 8.42$ Å, $Z = 4$. The structure was deduced from the three-dimensional Patterson synthesis and refined by block-diagonal least squares to a final agreement residual of 0.057. The molecule has nearly *mm* symmetry, and is folded to a dihedral form so that the angle between the slightly boat-shaped six-membered rings is 99.6°. The bond angle subtended by the oxygen atom is 101.4°, while the mean value of those in the methylene bridges is 119.0°. The strain implied by these values is consistent with the tendency of the compound to transform into the corresponding pyrene derivative, with ejection of the hetero atom.

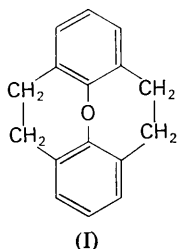
Introduction

Although several *trans*-[2.2]metacyclophanes have been reported (Lindsay, Stokes, Humber & Boekelheide, 1961), the title compound (I) is the first reported ex-

ample of a *cis*-[2.2]metacyclophane (Hess, Bailey & Boekelheide, 1967; Renfroe, Gurney & Hall, 1967). Steric strain is inferred for such compounds from their known tendency (on heating, or treatment with acid) to transform into the corresponding pyrene derivatives, with ejection of the hetero atoms (Boekelheide, 1967). The crystal structure analysis was undertaken in order to determine the geometrical characteristics of such strain. It was realized from the beginning that the ac-

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curacy of the analysis would be limited by intense thermal motion, with a consequent low yield of measurable reflections at high Bragg angles. Attempts to cool the specimen resulted in a phase change, with twinning. It was therefore necessary to accept the low accuracy inherent in the room-temperature determination.



Experimental

Crystal data: Orthorhombic

Formula $C_{16}H_{14}O$, F.W. 222.3

$a = 8.78 \pm 0.01$, $b = 16.04 \pm 0.02$, $c = 8.42 \pm 0.01$ Å

$Z = 4$, $V = 1185.8$ Å³

$D_x = 1.245 \pm 0.005$ g.cm⁻³, $D_m = 1.24 \pm 0.02$ g.cm⁻³

$\mu = 7.0$ cm⁻¹ (Cu $K\alpha$).

[Assumed wavelengths 1.54050 (Cu $K\alpha_1$); 1.54434 (Cu $K\alpha_2$)]

Space group $Cmc2_1$, $C2cm$, or $Cmcm$ (from precession photographs. Systematic absences: hkl for $h+k$ odd, $h0l$ for l odd). $Cmc2_1$ is consistent with the deduced structure.

Molecular symmetry: m .

The material supplied consisted of white translucent plates (010) of irregular outline. They appeared to possess surface charges, and readily adhered to probes. (This behavior suggested that the structure must lack a centre of symmetry.) The specimen used was a plate of approximate dimensions (initially) $0.05 \times 0.2 \times 0.3$ mm. During the eight days of data collection it suffered some decomposition, losing about half its original scattering power. The cell constants and relative intensities were measured with the General Electric XRD 5 spectrogoniometer and goniostat equipped with a scintillation counter. Copper $K\alpha$ radiation was used, and reasonable monochromatization was achieved with a $K\beta$ filter and a reverter (pulse-height analyser). The $\theta-2\theta$ scan method was used, with precomputed goniostat settings, and manual operation. Of the 744 accessible reflections (those for which $2\theta \leq 160^\circ$) significant counts were recorded for 448. The range of observed intensities was about 1 to 6000. Absorption corrections were considered to be unnecessary and were not made.

Structure determination

Packing considerations indicated that the probable space group was $Cmc2_1$, and a trial structure in this space group was deduced from the three-dimensional Patterson synthesis. Refinement was begun by means of Fourier syntheses, and completed with block-diagonal least-squares analysis. The program used was that of Ahmed, written in FORTRAN IV for the IBM system 360 computer. The quantity minimized is $\sum w(F_o - F_c)^2$. The matrices used are 3×3 for the posi-

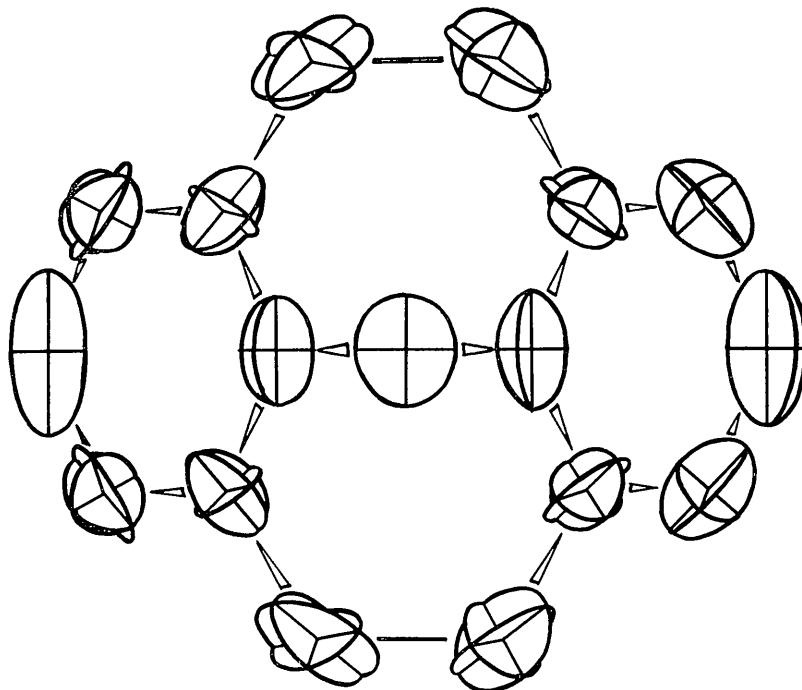


Fig. 1. Thermal ellipsoid plot of the molecule, viewed normal to its mean plane. 50% probability ellipsoids are represented.

tion parameters and 6×6 for the thermal parameters of each atom in a general position. The scale of F_o and the overall isotropic temperature factor are refined with a 2×2 matrix (Cruickshank, 1961). Schomaker's correction is applied to the shifts of the thermal parameters (Hodgson & Rollett, 1963). The weighting scheme used was $\sqrt{w} = 8/F_o$ for $F_o \geq 8$; $\sqrt{w} = F_o/8$ for $F_o < 8$. The nominal minimum observable value of F_o was 2.0. Refinement was rather slow, and it was necessary to prevent oscillation of parameters by applying a factor of 0.5 to all shifts. The scattering curves of Hanson, Herman, Lea & Skillman (1964) were used throughout. In the later stages of refinement the contribution of the hydrogen atoms was included in the structure factor calculations, assuming reasonable positions, and isotropic thermal motion. No attempt was made to refine the parameters of these atoms, however. At the same time a systematic extinction correction was derived in the manner of Pinnock, Taylor & Lipson (1956) and used to correct the intensities of the five strongest reflections. The greatest correction was about 40% of the uncorrected intensity. However, the weighting scheme used was such that these corrections were of little importance in the refinement. Unacceptable discrepancies between F_o and F_c were found for 9 weak reflections. These reflections (identified in Table 3) were omitted from the least-squares totals in the final cycles. Refinement was terminated when all shifts (after the damping factor of 0.5 had been applied) were less than one quarter of the corresponding e.s.d. The final parameters are given in Table 1.

Thermal motion

It is apparent from Table 1 that there is very intense thermal motion in the structure and that there must therefore be significant systematic errors in the observed bond lengths. Examination of the thermal ellipsoid plot (Fig. 1) suggests that at least some of the thermal motion occurs in rigid-body modes, and an analysis has therefore been carried out in the manner described by Cruickshank (1956). (It was assumed, without evidence, that the centre of libration coincided with the

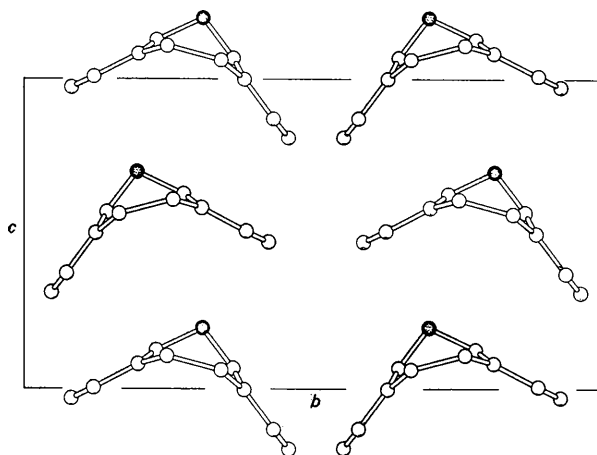


Fig. 2. The structure viewed along *a*. Light outlines represent molecules on the mirror plane at $x=0$; heavy represent those on the mirror plane at $x=\frac{1}{2}$.

Table 1(a). Final parameters (and e.s.d.'s) of non-hydrogen atoms

$$(0, 0, 0; \frac{1}{2}, \frac{1}{2}, 0) + (x, y, z; \bar{x}, y, z; x, \bar{y}, \frac{1}{2} + z; \bar{x}, \bar{y}, \frac{1}{2} + z)$$

$$\text{T.F.} = \exp \{-2\pi^2(U_{11}a^*2h^2 + 2U_{12}a^*b^*hk + 2U_{13}a^*c^*hl + U_{22}b^*2k^2 + 2U_{23}b^*c^*kl + U_{33}c^*2l^2)\}$$

(The numbering system is arbitrary. B_i are the principal-axis B_i 's.)

	x ($\times 10^4$)	y ($\times 10^4$)	z ($\times 10^4$)	U_{11} ($\times 10^4$)	U_{12} ($\times 10^4$)	U_{13} ($\times 10^4$)	U_{22} ($\times 10^4$)	U_{23} ($\times 10^4$)	U_{33} ($\times 10^4$)	B_1 (\AA^2)	B_2 (\AA^2)	B_3 (\AA^2)
O(1)	0	3063	1978	969	0	0	746	-66	487	7.65	6.01	3.72
		3	6	36			26	26	23			
C(2)	0	2272	1251	1115	0	0	475	120	604	8.80	5.34	3.18
		4	8	60			37	33	44			
C(3)	1389	1952	819	553	171	-206	698	-36	836	7.83	5.62	3.02
	5	3	6	27	29	35	26	26	31			
C(4)	1363	1195	74	879	287	-114	768	-33	1028	9.40	7.56	4.16
	6	3	8	37	35	42	27	30	37			
C(5)	0	810	-282	1887	0	0	609	-61	855	14.90	6.86	4.70
		4	10	94			43	50	61			
C(6)	0	3589	657	942	0	0	421	-188	541	7.44	5.36	2.24
		3	7	51			33	32	41			
C(7)	1385	3776	-24	654	-190	-67	592	-74	780	6.44	6.36	3.19
	5	3	6	29	29	35	21	24	29			
C(8)	1358	4282	-1319	877	-231	82	554	-30	1019	8.65	7.27	3.42
	6	3	7	34	27	38	25	29	38			
C(9)	0	4541	-1945	2087	0	0	419	90	747	16.48	6.08	3.12
		4	10	99			37	38	54			
C(10)	2834	2450	1076	800	221	-279	1015	-96	1184	11.24	7.76	4.68
	6	4	8	38	36	36	35	39	44			
C(11)	2829	3378	613	951	-257	-181	1195	-68	1219	10.78	10.07	5.72
	7	4	8	43	41	48	43	42	49			

centre of gravity of the equally-weighted non-hydrogen atoms.) The results are given in Table 2. It is clear from the smallness of the cross term in the tensor that the principal axes of libration are very close to the crystallographic axes. The axis of greatest libration is almost parallel to *b*, and thus about 13° from the axis of least moment of inertia. This seems physically reasonable; however, the agreement between observed and calculated U_{ij} 's is not good. The r.m.s. deviation is 0.0145 Å², compared with a mean e.s.d. for the observed elements of 0.0040 Å², and thus the hypothesis of pure rigid-body motion is not well supported. The reason may be systematic errors in the thermal parameters, or, quite probably, the presence of non-rigid modes of motion. In any case the bond lengths have not been corrected for rigid-body libration. The corrections would only be of the order of one e.s.d., and it is considered that errors resulting from non-rigid motion could be much greater than this.

Table 1(b). Assumed parameters for hydrogen atoms

	<i>x</i> (× 10 ⁴)	<i>y</i> (× 10 ⁴)	<i>z</i> (× 10 ⁴)	<i>B</i> (Å ²)
H(4)	2258	920	-50	4.4
H(5)	0	370	-908	5.4
H(8)	2305	4470	-1692	4.4
H(9)	0	4860	-2850	5.4
H(10 <i>a</i>)	3702	2176	398	5.9
H(10 <i>b</i>)	3142	2427	2245	5.9
H(11 <i>a</i>)	3602	3506	-242	5.9
H(11 <i>b</i>)	3078	3746	1570	5.9

Assessment of results

The agreement between F_o and F_c (Table 3) is reasonable, and there can be no doubt of the essential correctness of the structure. A difference Fourier synthesis of the cell contents projected along *c* reveals no detail inconsistent with the proposed structure. The agreement residual ($R = \Sigma ||F_o| - |F_c|| / \Sigma |F_o|$) is 0.057, for

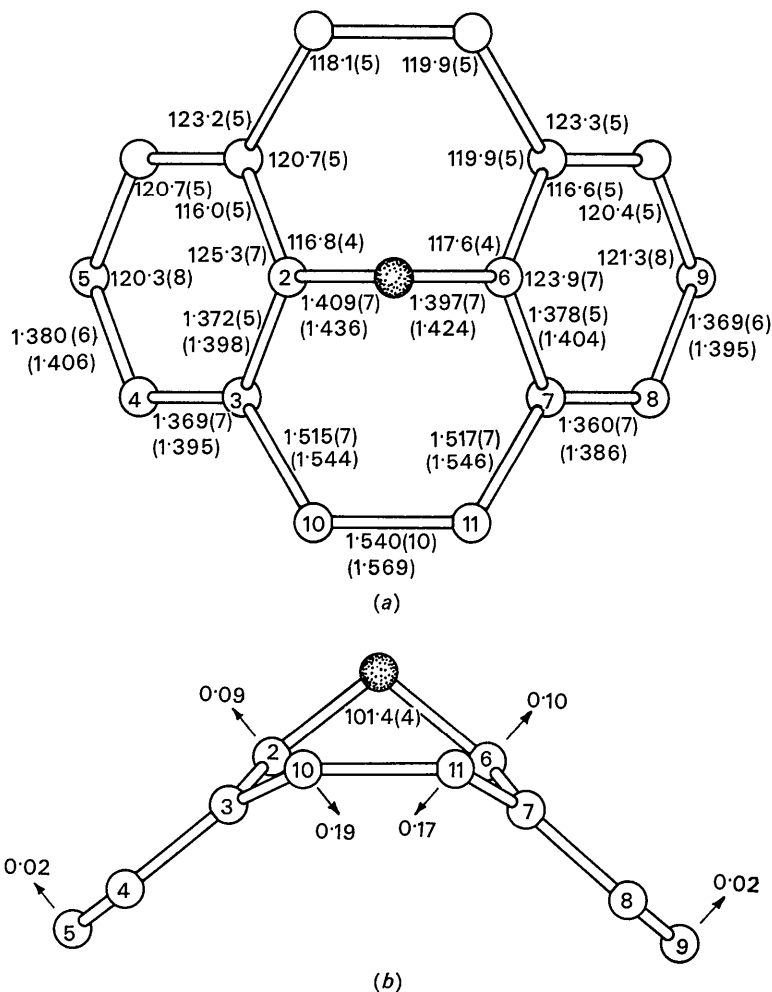


Fig. 3. Molecular geometry: Bond lengths and angles and their e.s.d.'s. The bond lengths in parentheses have been modified as described in the text. Arrows indicate the deviations (Å) of certain atoms from the planes defined by C(3)–C(4) and C(7)–C(8).

Discussion

The packing in the structure is illustrated in Fig. 2. The molecules are wedge-shaped, and all point in approximately the same direction. There could therefore be a permanent dipole associated with the crystal (a pyroelectric effect is allowed by the symmetry). The intermolecular distances have normal van der Waals values: there are fourteen such distances less than 4.00 Å, but none less than 3.62 Å.

The molecular geometry is illustrated in Fig. 3. The angle between the six-membered rings is 99.6°, and each is slightly distorted to a boat shape, as shown. There is no clear evidence of strain in the uncorrected lengths of the bridging C–O and C–C bonds. The CH₂–CH₂ distance of 1.540 Å does not differ significantly from the value of 1.533 Å reported for n-butane and n-hexane (Sutton, 1965). However, appropriately corrected distances (assuming the corrections could be estimated) might well give a different impression. An idea of the magnitude of such corrections is given by a comparison of the bond lengths in the six-membered rings of this and other structures. The mean value for benzene is 1.395 Å (*International Tables for X-ray Crystallography*, 1962) and for *trans*-4,12-dimethyl[2.2]-metacyclophane it is 1.399 Å (Hanson, 1962). The corresponding value in the present compound is 1.371 Å. There is no reason to expect such a reduction in the actual distance, and we might reasonably conclude that the observed distance should be increased by 0.026 Å, or 1.9%. If the remaining distances are increased in like manner (and this is clearly a dangerous procedure) the modified values are consistent with significant strain. In particular, the distance CH₂–CH₂ becomes 1.569 Å, a value which is comparable to the 1.573 Å found for *trans*-4,12-dimethyl[2.2]-metacyclophane. Although this result confirms reasonable expectations, it must be remembered that the modified distances are not very reliable, and do not by themselves constitute conclusive evidence of strain. Stronger evidence is provided by the differences of some of the bond angles from typical values. The angle C–O–C is 101.4°, quite different from the value of 111.5° found

in dimethyl ether (Sutton, 1965). The average value of the angles terminating the CH₂–CH₂ bond is 119.0°, whereas the unstrained tetrahedral value is 109.5°, and the corresponding value in the moderately-strained *trans*-4,12-dimethyl[2.2]-metacyclophane is only 111.5°. The deformation of this angle must provide at least some of the driving force for the ejection of the hetero atom in such compounds.

The problem was suggested, and the specimen material provided, by Professor V. Boekelheide. The computer programs used in the analysis are those of Ahmed, Hall, Pippy & Huber (1966). The assistance of these people, and the continued encouragement of Dr W. H. Barnes, are gratefully acknowledged.

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